# Surface Oxide Films and H<sub>2</sub>/CO Chemisorption at the Ru/TiO<sub>2</sub> Interface: Studies with a Model Planar Catalyst

JAS PAL S. BADYAL\* AND RICHARD M. LAMBERT<sup>+,1</sup>

\*Department of Chemistry, Science Laboratories, University of Durham, Durham DH1 3LE, United Kingdom; and †Department of Chemistry, University of Cambridge, Lensfield Road, Cambridge CB2 IEW, United Kingdom

Received July 2, 1990; revised January 25, 1991

The properties of high temperature treated (HTT) films of titanium oxide on Ru(0001) have been examined by electron spectroscopy, LEED, and CO/H<sub>2</sub> chemisorption. Starting with as-deposited TiO<sub>x</sub> films, surface-to-bulk transport begins at ~700 K provided that the initial oxide loading exceeds 0.68 ML. The system stabilises at ~1400 K with formation of a tenaciously held residual oxide film which is shown by LEED, Auger, and XPS measurements to be an expanded monolayer of TiO. No surface  $\rightarrow$  bulk transport occurs for TiO<sub>x</sub> loadings below the critical value of 0.68 ML. Chemisorptive properties have been studied as a function of HTT-TiO<sub>x</sub> film coverage and the results are discussed with respect to the SMSI phenomenon. (\* 1991 Academic Press, Inc.

### 1. INTRODUCTION

Planar model catalysts prepared by depositing support material on macroscopic metal specimens have been successfully employed to investigate aspects of metal-support interaction (1-3), and in the case of the SMSI phenomenon this approach can represent a particularly attractive strategy (4). It is well known that a number of different explanations has been advanced to explain SMSI behaviour, including migration of partially reduced oxide support species through or across the metal particles. Such behaviour would obviously strongly affect the availability of chemisorption sites and possibly also modify their reactive behaviour (5-8). An alternative type of explanation involves alloy formation between metal and support species (9-13) although the nature and constituent elements involved in such alloys are not clear.

In our earlier study of the  $Ru(0001)/TiO_2$ system, we showed how the presence of  $TiO_2$  moieties led to no more than simple site-blocking towards subsequent CO and  $H_2$  chemisorption (14). The present paper describes the effect of temperature on this system with specific reference to the formation of TiO<sub>x</sub> species that are stable at high temperature, their partitioning between surface and bulk, and the consequences of this behaviour for hydrogen and CO chemisorption. The TiO<sub>y</sub> films are produced by high temperature treatment and differ significantly in their surface chemistry from the unannealed TiO<sub>x</sub> (x > 2) films described in our earlier work (14). Use of a single-crystal specimen avoids the problems associated with grain boundary diffusion (15) which can dominate the behaviour of polycrystalline samples in experiments of this kind; we therefore have the possibility of determining unambiguously whether TiO, can be transported into the metal phase at temperatures relevant to SMSI (16).

#### 2. EXPERIMENTAL

A description of the apparatus and methods of specimen preparation and cleaning has been published elsewhere (17). Deposition of titanium oxide onto the single-crystal

<sup>&</sup>lt;sup>1</sup> To whom correspondence should be addressed.



FIG. 1. Variation of AES intensities with temperature for (a) 2.2 ML TiO<sub>x</sub>/Ru(0001) and (b) 0.6 ML TiO<sub>x</sub>/Ru(0001).

surface was carried out by controlled evaporation of Ti metal from a collimated source operating in an ambient atmosphere of  $1 \times$  $10^{-6}$  Torr O<sub>2</sub>. We have shown previously (14) that this procedure results in the formation of an ultrathin layer of  $TiO_x$  film whose as-deposited stoichiometry corresponds to x = 2 at the point of monolayer completion. Selective removal of oxygen chemisorbed on the exposed Ru metal surface is achieved by holding the specimen at 573 K and irradiating with a flux of hydrogen atoms and ions issuing from a hot cathode low energy source (14), a procedure which does not significantly perturb the TiO<sub>x</sub> film itself. By this means, patches of TiO<sub>y</sub> were generated on an otherwise bare Ru surface. In obtaining the results which follow, each experiment was carried out by rigorously cleaning the surface (17), depositing a fresh dose of  $TiO_x$ , followed by selective removal of the oxygen chemisorbed on Ru. Quoted TiO<sub>y</sub> coverages were obtained by means of a calibration previously established with the aid of LEED/ Auger spectroscopy data (14); 1.0 monolayer (ML) refers to a monolayer film of asdeposited TiO<sub>x</sub> at room temperature.

### 3. RESULTS

## 3.1 Temperature-Dependent Behaviour of TiO<sub>x</sub> on Ru(0001)

Figure 1A shows the behaviour of an initial 2.2-monolayer film of  $TiO_x$  as a function of increasing temperature. The data were obtained by following the Auger electron signal intensities due to Ru (231 eV), Ti (387 eV), and O (510 eV), in each case after heating the sample for a fixed interval of 60 s at various increasing temperatures. In any given case, it was found that quenching the sample to room temperature led to no observable change in signal intensities, so that the data were acquired at room temperature. It can be seen that at sufficiently high heating temperatures (~1350 K) all three signals tend to limiting values. This type of behaviour could be due either to agglomeration of TiO<sub>x</sub> at the surface or to diffusion of either Ti or Ti + O into the bulk of the metal.

Once the system had reached its limiting high temperature condition, limited Ar<sup>+</sup> etching of the surface at 300 K resulted in an Auger spectrum corresponding to clean Ru, essentially all Ti and oxygen having been removed. This is illustrated in Fig. 2 which shows a sequence of Auger spectra corresponding to (A) an initially clean surface, (B) after depositing 1.1 ML at room temperature, (C) the result of subsequent heating to 1550 K, and (D) following Ar<sup>+</sup> cleaning at 300 K followed by annealing to 600 K. Further annealing to 1550 K followed by cooling to room temperature led to the appearance of Auger signals characteristic of TiO<sub>r</sub> (Fig. 2E). Heating in oxygen (120 L at 1300 K) indicated no further segregation of  $TiO_x$  (Fig. 2F). (This procedure is known to leach out dissolved titanium from bulk ruthenium (17)). It is noteworthy that, as judged by Auger intensities, the



FIG. 2. (A) AES spectrum of clean Ru(0001); (B) after deposition of  $1.1 \text{ ML TiO}_x$  at 295 K; (C) following a flash to 1550 K; (D) argon ion cleaned and annealed to 600 K; (E) following flash to 1550 K; (F) after heating in 120 L O<sub>2</sub> at 1300 K followed by flashing to 1550 K.

amount of  $\text{TiO}_x$  initially present (i.e., before annealing) is equal to the amount remaining after annealing plus the amount dissolved (and subsequently re-exposed by etching and segregation).

These data imply that high temperature treatment of TiO<sub>x</sub> overlayers on ruthenium result in a limiting coverage of TiO<sub>x</sub> (0.68  $\pm$  0.04 ML) which remains on the surface while the remainder diffuses into the bulk. For initial loadings of TiO<sub>x</sub> less than this critical amount, no diffusion of TiO<sub>x</sub> into the bulk occurs (Fig. 1B). Equally, for initial loadings of greater than 0.68 ML, annealing to ~1500 K results in surface-to-bulk diffusion until a coverage of 0.68 ML remains at the surface (Fig. 3). Figure 4A shows the LEED pattern characteristic of this 0.68-ML TiO<sub>x</sub> overlayer. It may be interpreted in terms of a  $(5\sqrt{3} \times 5\sqrt{3})R30^{\circ}$  coinci-

dence lattice formed between the TiO<sub>x</sub> overlayer and the Ru(0001) substrate (Fig. 4B). In terms of this model, the calculated fractional coverage TiO, /Ru is 0.64  $\pm$  0.04, in very good agreement with the overlayer coverage as determined by Auger spectroscopy. XP spectra (Fig. 5) taken from freshly deposited TiO, overlayers and from the annealed system revealed substantial shifts in titanium core level binding energies (BEs). For the as-deposited TiO<sub>y</sub> film the Ti $(2p_{3/2})$ BE was 459.2 eV. Annealing of such films to ~1500 K to produce the 0.68-ML TiO<sub>x</sub> phase gave a substantially reduced  $Ti(2p_{3/2})$ BE of 456.0 eV. This 3.2 eV decrease in BE is consistent with the transformation TiO<sub>2</sub>  $\rightarrow$  TiO (18). Desorption measurements were carried out on multilayer as-deposited films (i.e., unannealed) using a collimated quadrupole mass spectrometer with good detection sensitivity for species desorbing from the front face of the specimen. No signals were detected at mass numbers of  $32(O_2)$ , 48(Ti), 64(TiO), or  $80(TiO_2)$ ; this tends to confirm that the changes in TiO<sub>x</sub> coverage which result from heating are due to surface  $\rightarrow$  bulk diffusion and not the evaporation of surface species.

# 3.2 Carbon Monoxide and Hydrogen Chemisorption

The known TPD fingerprint spectrum for CO from clean Ru(0001) is used to investigate the influence of the reduced titanium oxide species (i.e., after annealing) on



FIG. 3. AES signals after annealing to 1550 K as a function of TiO<sub>v</sub>/Ru(0001) precoverage.



FIG. 4. (A) LEED pattern of TiO<sub>x</sub>/Ru(0001); 1.1 ML precoverage flashed to 1550 K (beam energy = 131.7 eV); (B) showing  $(5\sqrt{3} \times 5\sqrt{3}))R30^\circ$  unit mesh of corresponding coincidence lattice.

neighboring ruthenium sites, and the results are shown in Fig. 6A for a range of  $\text{TiO}_x$ precoverages. It can be seen that with increasing  $\text{TiO}_x$  coverage, the CO desorption shifts to higher temperatures. Furthermore, examination of the CO desorption yield as a function of  $\text{TiO}_x$  coverage (Fig. 6B) shows a non-linear attenuation of CO uptake; this is interesting because the behaviour is markedly different from the simple linear CO site blocking behaviour exhibited by unreduced titanium oxide species (i.e., un-annealed oxide films) (14). (In Fig. 6B the uptake at zero  $\text{TiO}_x$  coverage corresponds to saturation coverage of CO on bare Ru(0001).)

The essentially complete suppression of CO uptake at TiO<sub>x</sub> coverages below 1 ML (Fig. 6B) is consistent with expansion of the TiO<sub>x</sub> overlayer upon annealing *for room temperature precoverages* < 0.64 ML TIO<sub>x</sub>. Adsorption/desorption experiments carried out with <sup>13</sup>C<sup>16</sup>O/<sup>12</sup>C<sup>18</sup>O mixtures showed no evidence for isotopic scrambling indicating that only molecular CO adsorbs and desorbs from the TiO<sub>x</sub>-modified Ru surface.

In our earlier work (14) we reported that hydrogen chemisorption was drastically inhibited by as-deposited  $\text{TiO}_x$  (i.e., unannealed films) and that subsequent desorption of H<sub>2</sub> was shifted to lower temperature. In the present case, the decrease in H<sub>2</sub> chemisorption with *annealed*  $\text{TiO}_x$  precoverage is less marked, taking account of the fact that in this case 0.64 ML of  $\text{TiO}_x$  precoverage is sufficient to completely block the ruthenium surface (Figs. 7A,7B). Furthermore, in this case, as distinct from the asdeposited case, there is a shift of hydrogen desorption to *higher* temperatures with increasing  $\text{TiO}_x$  precoverage.

### 4. DISCUSSION

The present results show that the method and temperature of preparation of  $\text{TiO}_x$ overlayers on metal surfaces can be critical in determining the metal coverage and chemisorption behaviour of such systems. Such factors may have led to apparent discrepancies resulting in disagreements about the actual role of TiO<sub>x</sub> moieties on



FIG. 5. XP spectra as a function of annealing temperature of a TiO<sub>x</sub> film. Note that the large feature around 462 eV is mainly due to the Ru( $3p_{3/2}$ ) level with some contribution from Ti( $2p_{1/2}$ ).

metal particles (19). Thus in some cases it has been argued that the effect of such species is simply to block metal sites to chemisorption (20, 21), whereas other reports indicate that a longer range effect operates, with adjacent metal sites also being affected by adsorbed  $\text{TiO}_x$  (22, 23).

The formation of a residual tenaciously held TiO<sub>y</sub> overlayer after high temperature treatment is somewhat reminiscent of the behaviour exhibited by Ti films on Ru (24): however, in the present case, the limiting Ti coverage (0.68 ML) is greater than with the metal-only system (0.49 ML). This presumably reflects the effect of oxygen-Ti binding in stabilising Ti at the metal/gas interface. Not surprisingly perhaps, transport of TiO<sub>y</sub> into Ru is also significantly harder than that of Ti alone; diffusion becomes appreciable at  $\sim$ 700 K in the former case and at  $\sim$ 500 K in the latter. It is not possible to comment on whether in the present case the diffusion mechanism involves transport of an intact titanium oxide species as opposed to decomposition



FIG. 6. (A) TPD spectra as a function of  $TiO_x$  precoverage following 44-L CO doses at 140 K; (B) molecular CO uptake per surface ruthenium atom.



FIG. 7. (A) TPD spectra as a function of TiO<sub>x</sub> precoverage following 200-L H<sub>2</sub> doses; (B) amount of H uptake per surface ruthenium atom.

of TiO<sub>x</sub> at the surface followed by substitutional diffusion of Ti and interstitial diffusion of the associated oxygen atom. What is clear, however, is that the presence of Ti substantially lowers the activation energy for oxygen diffusion into bulk ruthenium because it is well known that oxygen alone does not diffuse into Ru metal under conditions similar to those implied here (25).

Taken together, the Auger, XPS, LEED, and TPD observations are consistent with the view that high temperature treatment of  $TiO_x/Ru(0001)$  leads to a reduction in oxidation state of the  $Ti^{n+}$  ions. (Our earlier work on as-deposited submonolayers of  $TiO_x$  indicated n > 2 for such systems (14).) The lattice parameter of defective titanium monoxides is proportional to the degree of reduction (26) and this is in good agreement with the observed expansion of the  $TiO_x$ monolayer film deduced from the Auger and LEED data; the XPS observations confirm a reduction in Ti oxidation state as a result of high temperature treatment.

It seems plausible that the presence of reduced Ti ions in the high temperature treated TiO<sub>y</sub> film should lead to some electronic charge transfer to neighbouring Ru atoms resulting in the observed increase in CO chemisorption strength on these atoms. On clean Ru(0001) CO molecules are linearly bonded at on-top sites, and at high coverages interactions in the compressed adsorbate layer are accommodated by a marked tilt of the CO molecules away from surface normal (27, 28). The non-linear decrease in CO uptake with increasing coverage of high temperature treated TiO<sub>y</sub> may reflect the increased Ru-CO interaction which could inhibit the formation of a compressed CO overlayer.

The corresponding behaviour towards hydrogen chemisorption may be understood as follows. For as-deposited  $\text{TiO}_x$  films (14) highly dispersed titanium oxide species decrease the number of Ru ensembles available for dissociative chemisorption of H<sub>2</sub>; high temperature treatment leads to the coalescence of the  $\text{TiO}_x$  species to produce larger islands (detectable by LEED) leaving a larger number of active ensembles and increasing the hydrogen uptake capacity. An additional possibility is that edges of the HTT-TiO<sub>x</sub> islands function as active centres for H<sub>2</sub> dissociation (29) followed by spillover onto neighbouring Ru. It is interesting to note that in CO/H<sub>2</sub> competitive chemisorption, any changes which favour the relative amount of H<sub>ad</sub> could enhance the propagation step in Fischer–Tropsch synthesis, thereby affecting the product distribution (30–32), as has indeed been reported for systems exhibiting the SMSI effect (33).

### 5. CONCLUSIONS

The surface properties and chemisorptive behaviour of TiO<sub>y</sub> on Ru are strongly dependent on temperature and initial loading. The LEED and Auger data show that high temperature treatment can result in material transport and formation of a new surface phase. In particular, if the oxide loading exceeds a critical value, high temperature treatment results in surface-to-bulk diffusion of both titanium and oxygen leading ultimately to the formation of a stable residual monolayer film of expanded TiO. The XPS results show that reduced titanium species are formed after high temperature treatment. These reduced titanium species at the metal surface transfer electronic charge to Ru; TPD observations reveal associated changes in CO and hydrogen chemisorption behaviour which closely resemble those reported for catalysts in the SMSI condition. It is noteworthy that these phenomena of oxide dissolution and residual TiO surface film formation occur at temperatures similar to those required to induce SMSI behaviour in practical catalysts.

### ACKNOWLEDGMENT

We are grateful to Johnson-Matthey plc for a loan of precious metals.

### REFERENCES

I. Strongin, D. R., Bare, S. R., and Somorjai, G. A., J. Catal. 103, 289 (1987).

- Demmin, R. A., and Gorte, R. J., J. Catal. 105, 373 (1987).
- Zhao, Y.-B., and Chung, Y.-W., J. Catal. 106, 369 (1987).
- Tauster, S. J., Fung, S. C., and Garten, R. L., J. Am. Chem. Soc. 100, 170 (1978).
- 5. Meriaudeau, P., Dutel, J. F., Dufaux, M., and Naccache, C., Stud. Surf. Sci. Catal. 11, 95 (1982).
- Santos, J., Phillips, J., and Dumesic, J. A., J. Catal. 81, 147 (1983).
- Resasco, D. E., and Haller, G. L., J. Catal. 82, 279 (1983).
- Sadeghi, H. R., and Henrich, V. E., J. Catal. 87, 279 (1984).
- Baker, R. T. K., Prestridge, E. B., and Garten, R. L., J. Catal. 56, 390 (1979).
- 10. Cairns, J. A., Baglin, J. E. E., Clark, G. J., and Ziegler, J. F., J. Catal. 83, 301 (1983).
- Sakellson, S., McMillan, M., and Haller, G. L., J. Phys. Chem. 90, 1733 (1986).
- 12. Spencer, M. S., J. Catal. 93, 216 (1985).
- Beard, B. C., and Ross, P. N., J. Phys. Chem. 90, 6811 (1986).
- 14. Badyal, J. P. S., Gellman, A. J., Judd, R. W., and Lambert, R. M., *in* "Structure and Reactivity of Surfaces" (C. Morterra, A. Zecchina, and G. Costa, Eds.), Stud. Surf. Sci. Catal., Vol. 48, pp. 19–30. Elsevier, Amsterdam, 1988.
- Reed-Hill, R. E., "Physical Metallurgy Principles," 2nd ed., p. 418. Brooks/Cole Engineering Division, 1973.
- 16. Ko, C. S., and Gorte, R. J., J. Catal. 90, 59 (1984).
- Badyal, J. P. S., Gellman, A. J., and Lambert, R. M., Surf. Sci. 188, 557 (1987).
- Carley, A. F., Chalker, P. R., Rivière, J. C., and Roberts, M. W., *J. Chem. Soc. Faraday Trans. 1* 83, 351 (1987).
- Gorte, R. J., Altman, E., Corallo, G. R., Davidson, M. R., Ashbury, D. A., and Horflund, G. B., *Surf. Sci.* 188, 327 (1987).
- Dwyer, D. J., Cameron, S. D., and Gland, J., Surf. Sci. 159, 430 (1985).
- 21. Ko, C. S., and Gorte, R. J., Surf. Sci. 161, 597 (1985).
- 22. Takatani, S. and Chung, Y. W., Appl. Surf. Sci. 19, 341 (1984).
- Levin, M., Salmeron, M., Bell, A. T., and Somorjai, G. A., Surf. Sci. 169, 123 (1986).
- 24. Badyal, J. P. S., Nix, R. M., and Lambert, R. M., *Faraday Discuss. Chem. Soc.* 87, 121 (1989).
- Praline, G., Koel, B. E., Lee, H.-I., and White, J. M., Appl. Surf. Sci. 5, 296 (1980).
- Donnay, J. D. H., Donnay, G., Cox, E. G., Kennard, O., and King, M. V., "Crystal Data Determinative Tables," 2nd ed. American Crystallographic Society, 1963.
- Treichler, R., Riedl, W., Wurth, W., Feulner, P. and Menzel, D., *Phys. Rev. Lett.* 54, 462 (1985).

- 28. Riedl, W., and Menzel, D., Surf. Sci. 163, 39 (1985).
- Badyal, J. P. S., Gellman, A. J., and Lambert, R. M., J. Catal. 111, 383 (1988).
- 30. Wang, S. Y., Moon, S. H., and Vannice, M. A., J. Catal. 56, 167 (1981).
- 31. Raupp, G. B., and Dumesic, J. A., J. Catal. 96, 597 (1985).
- Flory, P. J., "Principles of Polymer Chemistry,"
  p. 318. Cornell Univ. Press, Ithaca, New York, 1967.
- 33. Vannice, M. A., and Garten, R. L., J. Catal. 63, 255 (1980).